

# New directional signatures from the non-relativistic EFT of dark matter or *'Who ordered all these operators?'*

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RPP 2016, Annecy - 25th Jan. 2016

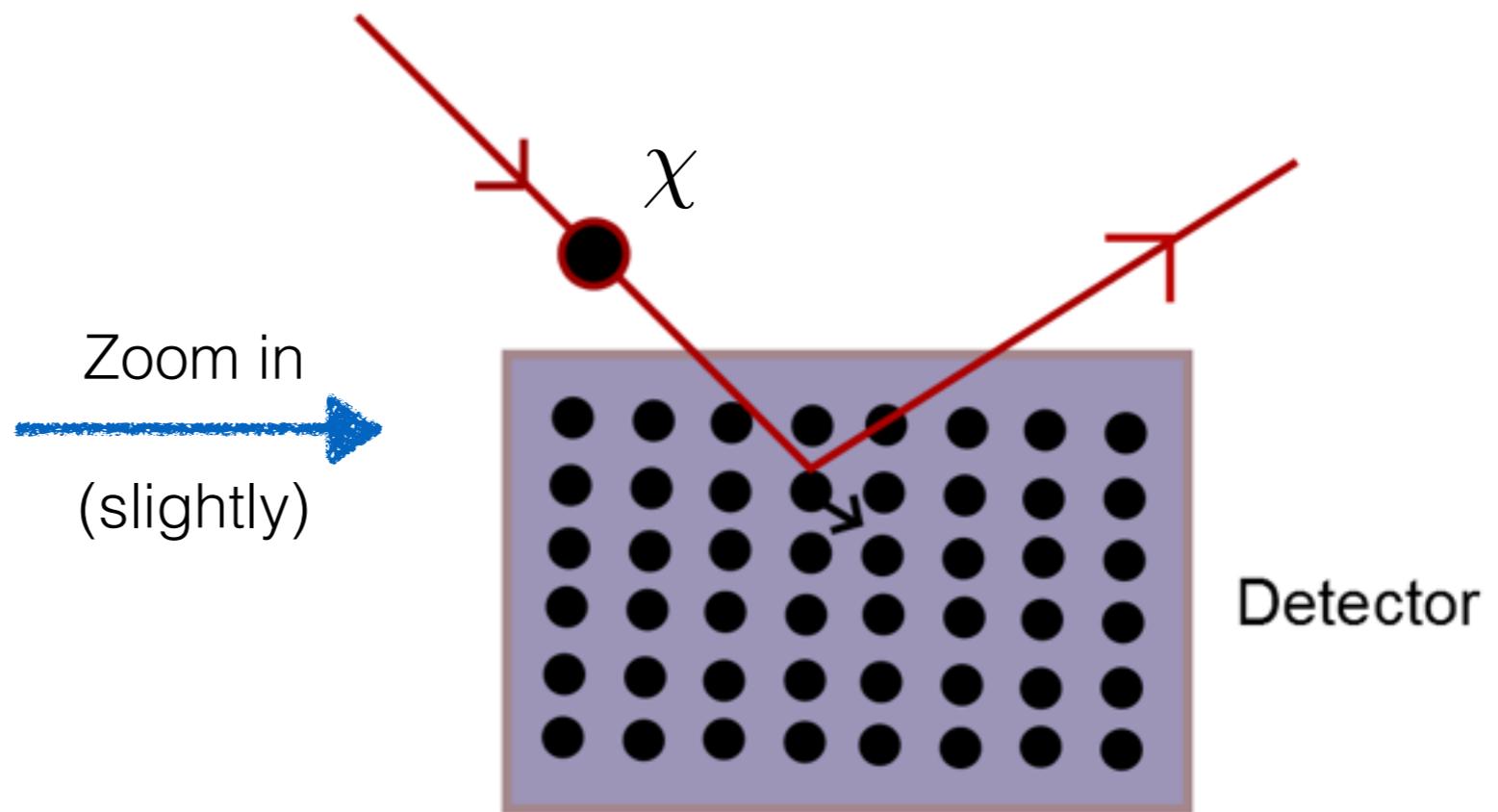
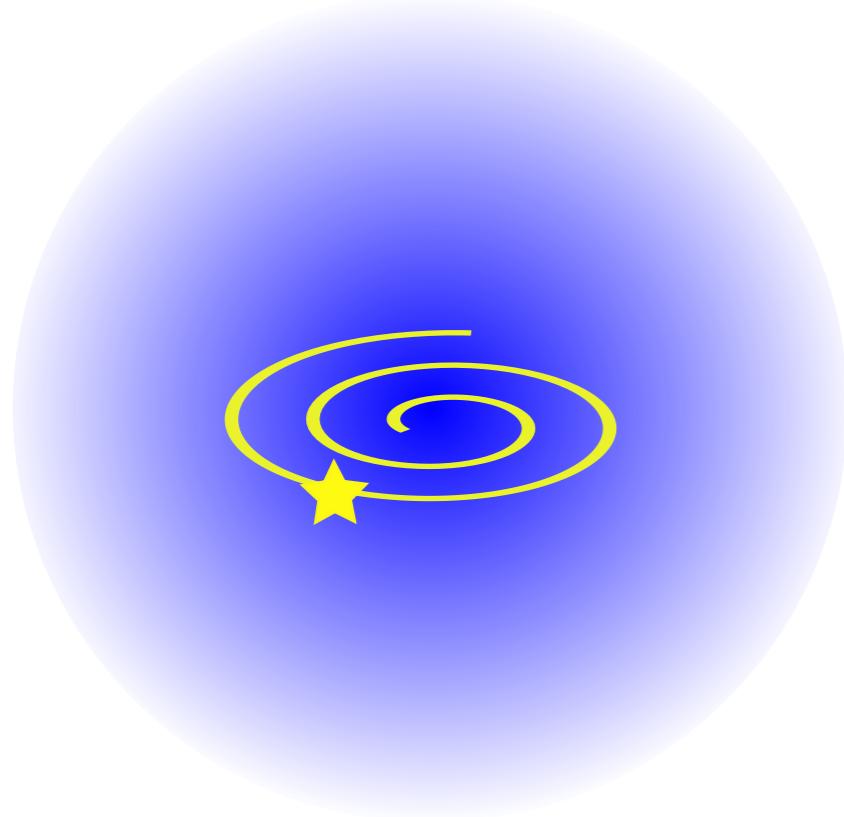


John  
Templeton  
Foundation

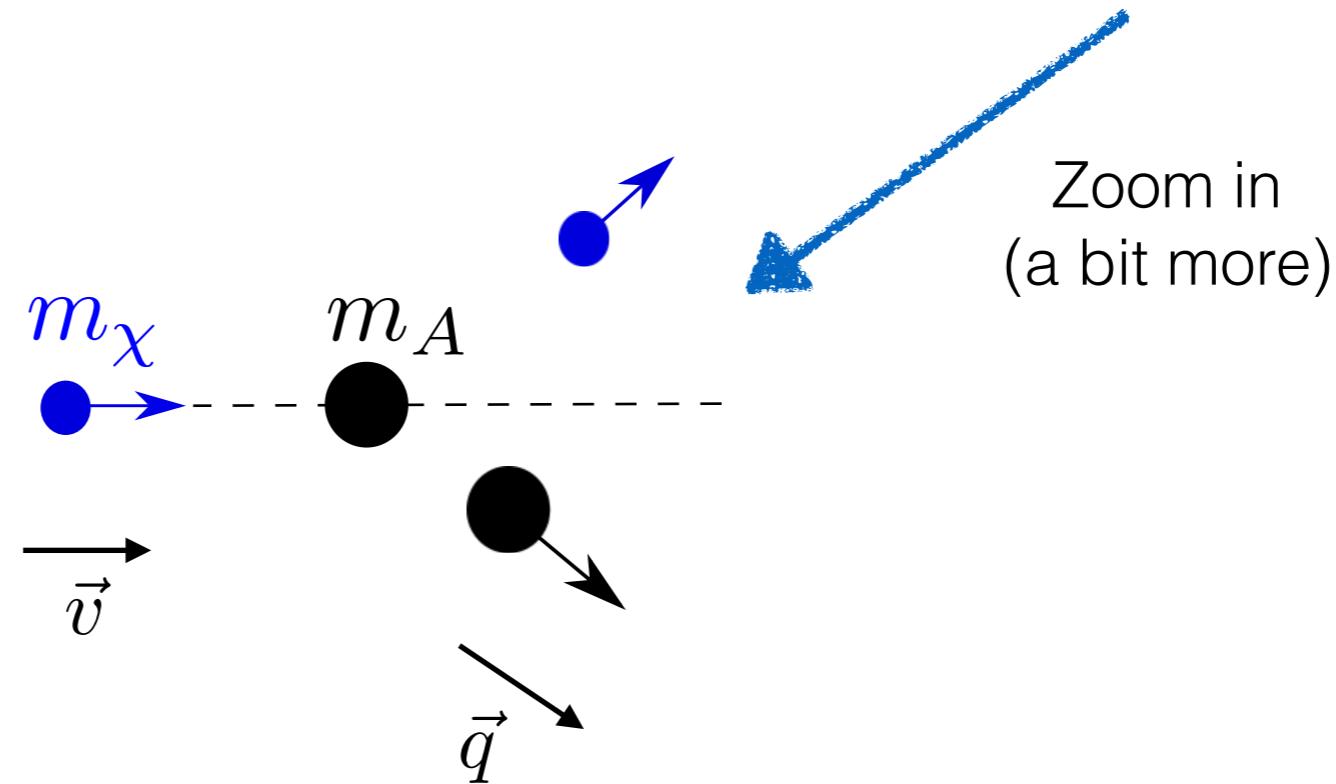


Based on  
arXiv:1505.07406

# Direct detection of Dark Matter



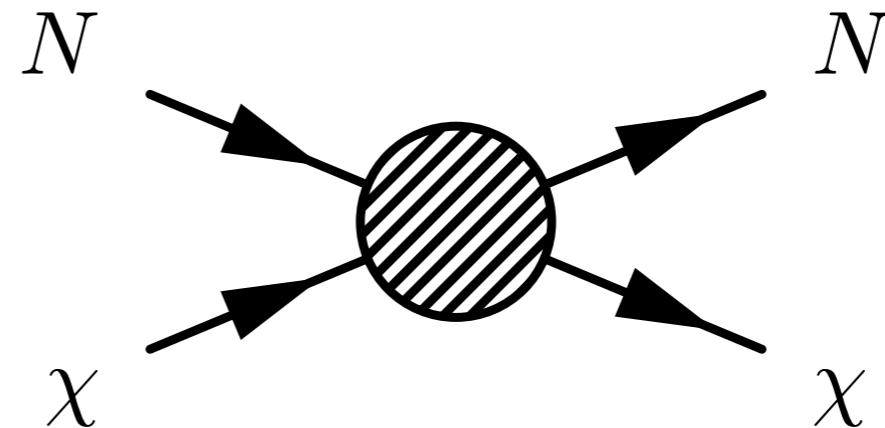
$$m_\chi \gtrsim 1 \text{ GeV}$$
$$v \sim 10^{-3}$$



# DM-nucleon interactions

Direct detection:

$$m_\chi \gtrsim 1 \text{ GeV}$$
$$v \sim 10^{-3}$$

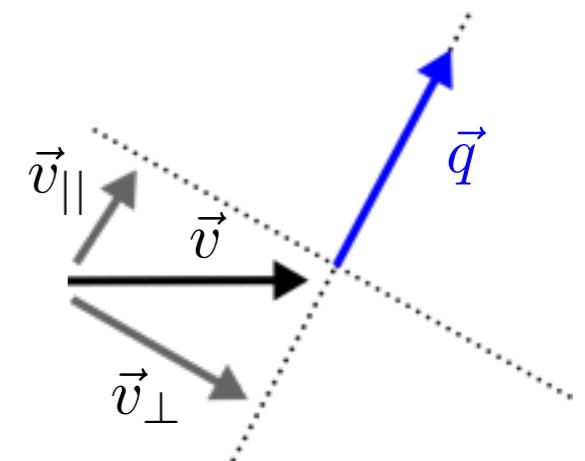


$$q \lesssim 100 \text{ MeV} \sim (2 \text{ fm})^{-1}$$

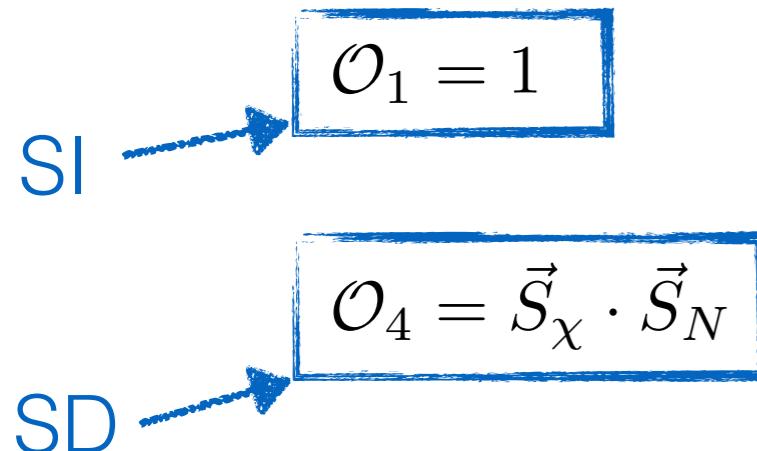
Relevant non-relativistic (NR) degrees of freedom:

$$\vec{S}_\chi, \quad \vec{S}_N, \quad \frac{\vec{q}}{m_N}, \quad \vec{v}_\perp = \vec{v} + \frac{\vec{q}}{2\mu_{\chi N}}$$

Fitzpatrick et al. [arXiv:1203.3542]



# Non-relativistic effective field theory (NREFT)



[arXiv:1008.1591, arXiv:1203.3542, arXiv:1308.6288, arXiv:1505.03117]

# Non-relativistic effective field theory (NREFT)

$\mathcal{O}_1 = 1$	$\mathcal{O}_{12} = \vec{S}_\chi \cdot (\vec{S}_N \times \vec{v}^\perp)$
$\mathcal{O}_3 = i\vec{S}_N \cdot (\vec{q} \times \vec{v}^\perp)/m_N$	$\mathcal{O}_{13} = i(\vec{S}_\chi \cdot \vec{v}^\perp)(\vec{S}_N \cdot \vec{q})/m_N$
$\mathcal{O}_4 = \vec{S}_\chi \cdot \vec{S}_N$	$\mathcal{O}_{14} = i(\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{v}^\perp)/m_N$
$\mathcal{O}_5 = i\vec{S}_\chi \cdot (\vec{q} \times \vec{v}^\perp)/m_N$	$\mathcal{O}_{15} = -(\vec{S}_\chi \cdot \vec{q})((\vec{S}_N \times \vec{v}^\perp) \cdot \vec{q}/m_N^2$
$\mathcal{O}_6 = (\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{q})/m_N^2$	⋮
$\mathcal{O}_7 = \vec{S}_N \cdot \vec{v}^\perp$	
$\mathcal{O}_8 = \vec{S}_\chi \cdot \vec{v}^\perp$	
$\mathcal{O}_9 = i\vec{S}_\chi \cdot (\vec{S}_N \times \vec{q})/m_N$	
$\mathcal{O}_{10} = i\vec{S}_N \cdot \vec{q}/m_N$	
$\mathcal{O}_{11} = i\vec{S}_\chi \cdot \vec{q}/m_N$	

SI →  $\mathcal{O}_1 = 1$  and  $\mathcal{O}_4 = \vec{S}_\chi \cdot \vec{S}_N$

SD →  $\mathcal{O}_5 = i\vec{S}_\chi \cdot (\vec{q} \times \vec{v}^\perp)/m_N$

[arXiv:1008.1591, arXiv:1203.3542, arXiv:1308.6288, arXiv:1505.03117]

# Non-relativistic effective field theory (NREFT)

SI

$$\mathcal{O}_1 = 1$$

$$\mathcal{O}_3 = i\vec{S}_N \cdot (\vec{q} \times \vec{v}^\perp)/m_N$$

$$\mathcal{O}_4 = \vec{S}_\chi \cdot \vec{S}_N$$

SD

$$\mathcal{O}_5 = i\vec{S}_\chi \cdot (\vec{q} \times \vec{v}^\perp)/m_N$$

$$\mathcal{O}_6 = (\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{q})/m_N^2$$

$$\mathcal{O}_7 = \vec{S}_N \cdot \vec{v}^\perp$$

$$\mathcal{O}_8 = \vec{S}_\chi \cdot \vec{v}^\perp$$

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$$\mathcal{O}_{14} = i(\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{v}^\perp)/m_N$$

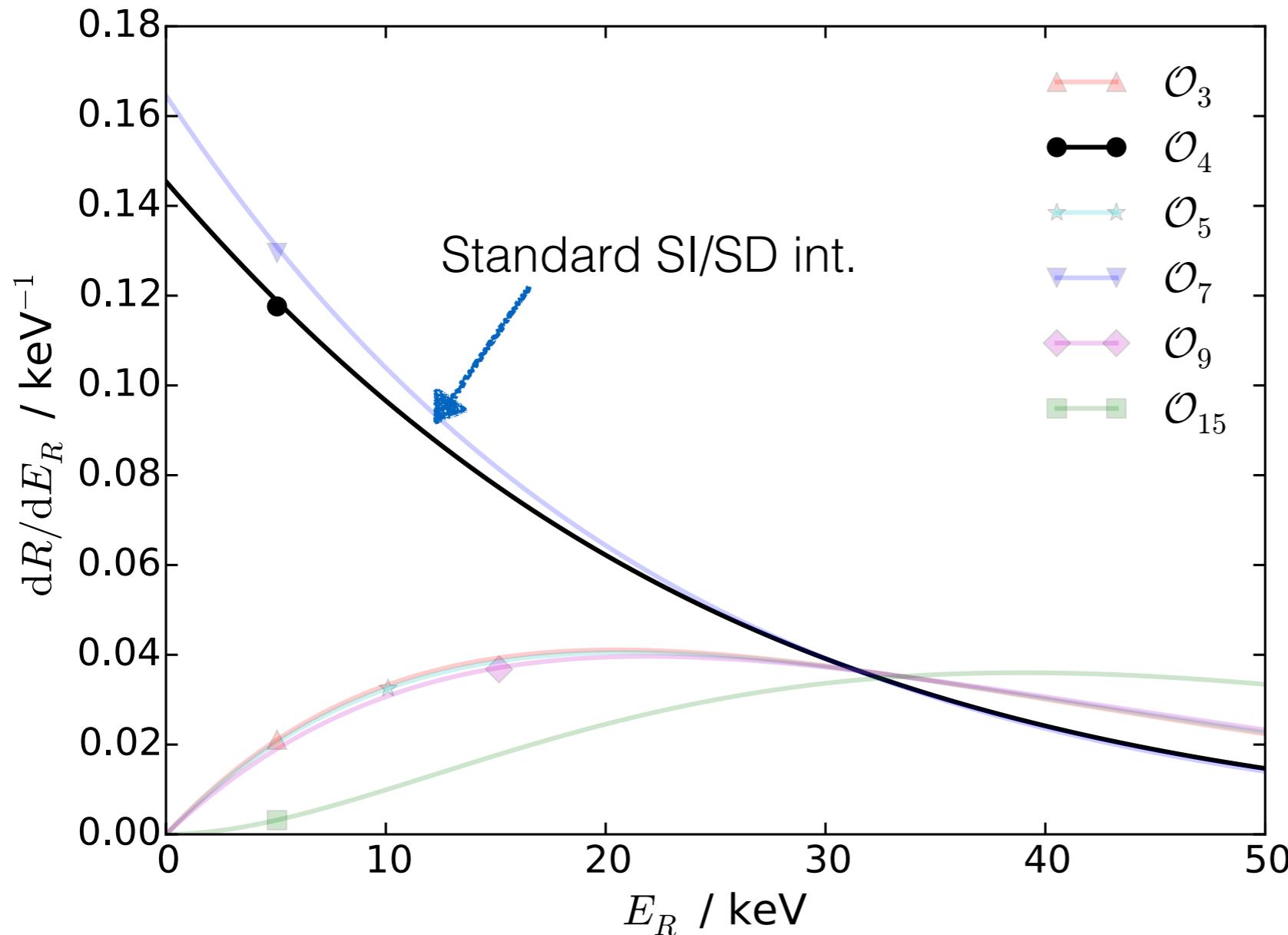
$$\mathcal{O}_{15} = -(\vec{S}_\chi \cdot \vec{q})((\vec{S}_N \times \vec{v}^\perp) \cdot \vec{q}/m_N^2$$

⋮

$$\frac{d\sigma_i}{dE_R} \sim \frac{1}{v^2} F_i(v_\perp^2, q^2)$$

[arXiv:1008.1591, arXiv:1203.3542, arXiv:1308.6288, arXiv:1505.03117]

# Standard energy spectrum



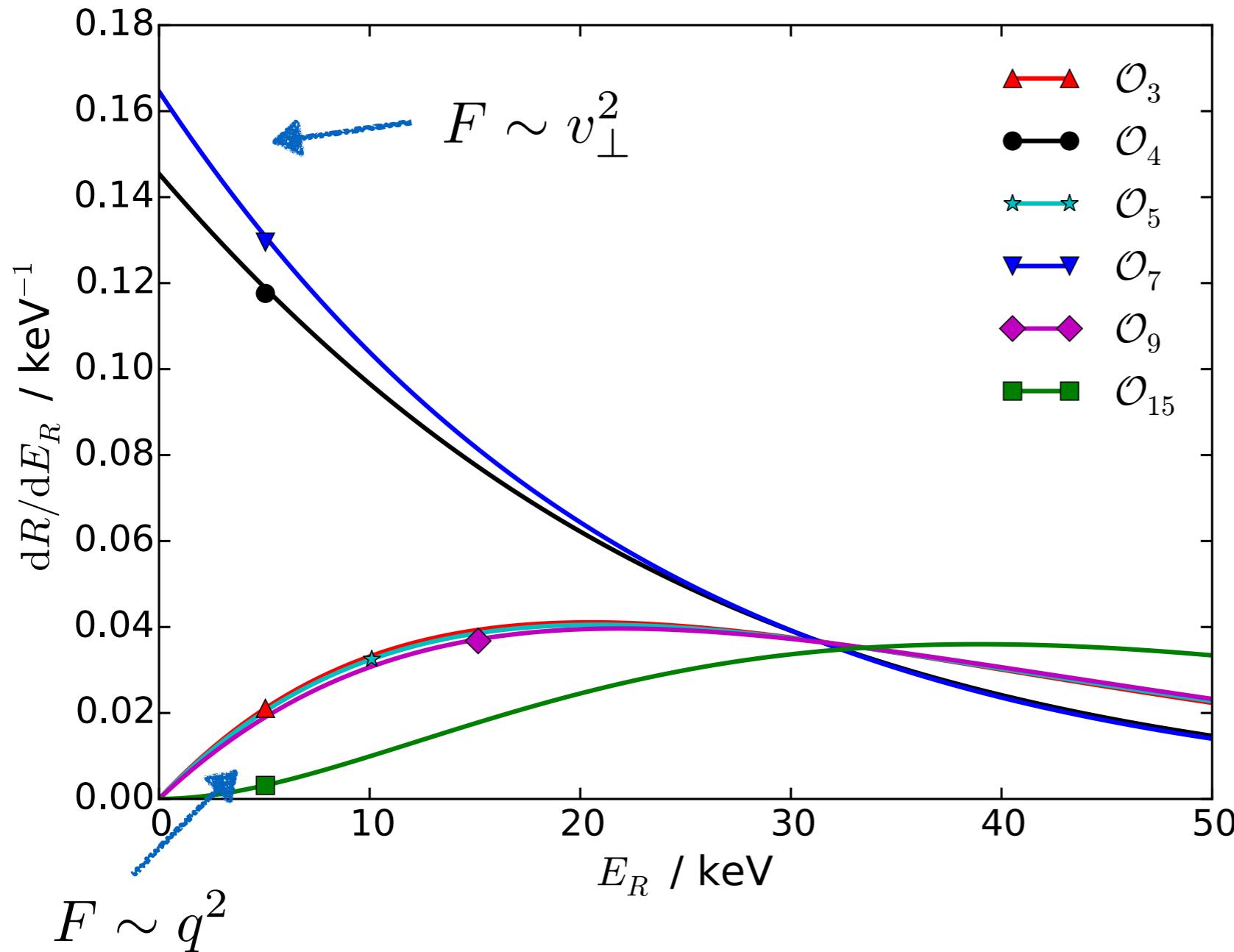
'Perfect' CF<sub>4</sub> detector

$E_R \in [20, 50] \text{ keV}$

Input WIMP mass:  
 $m_\chi = 100 \text{ GeV}$

SHM velocity distribution

# NREFT energy spectrum



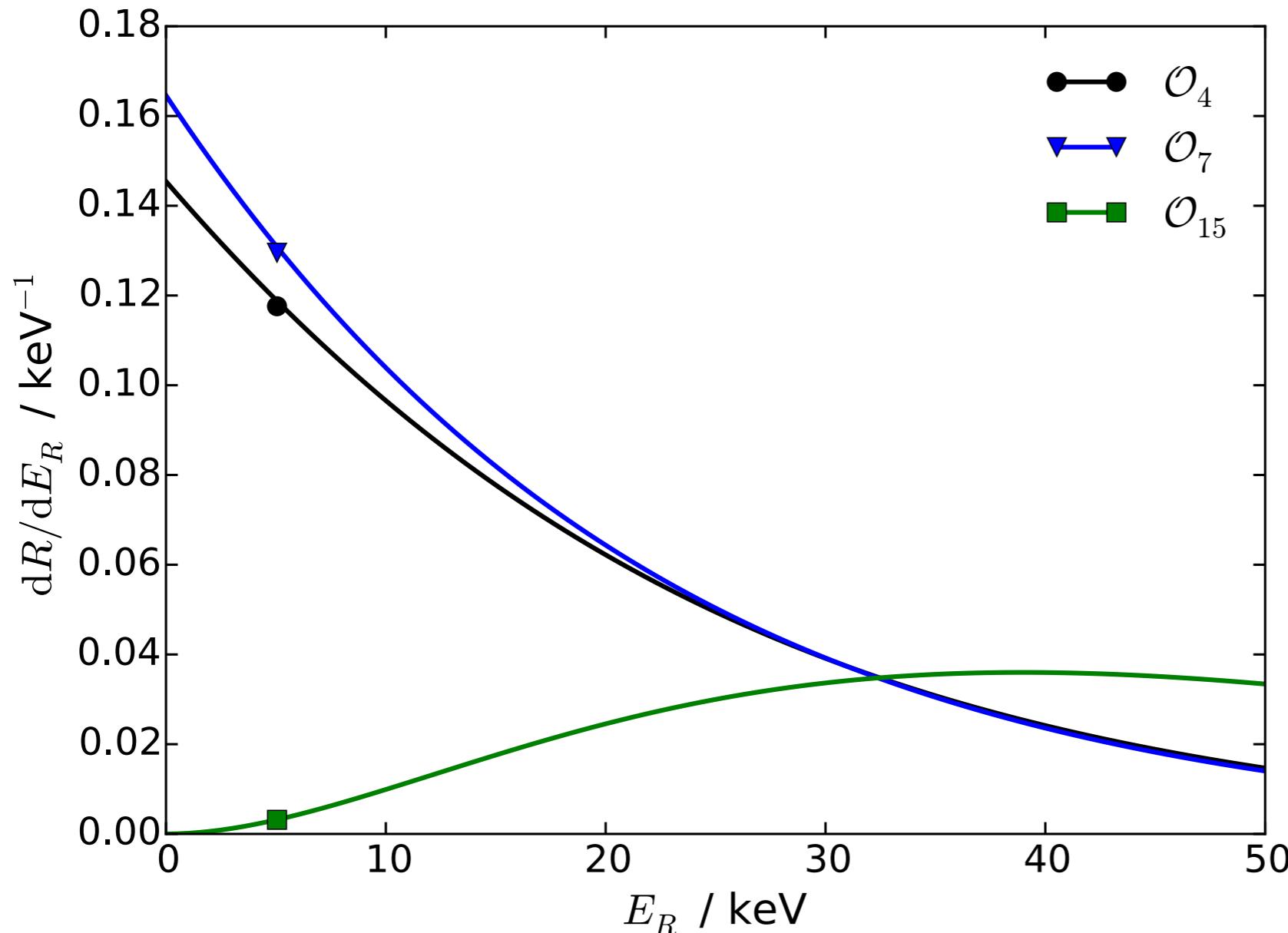
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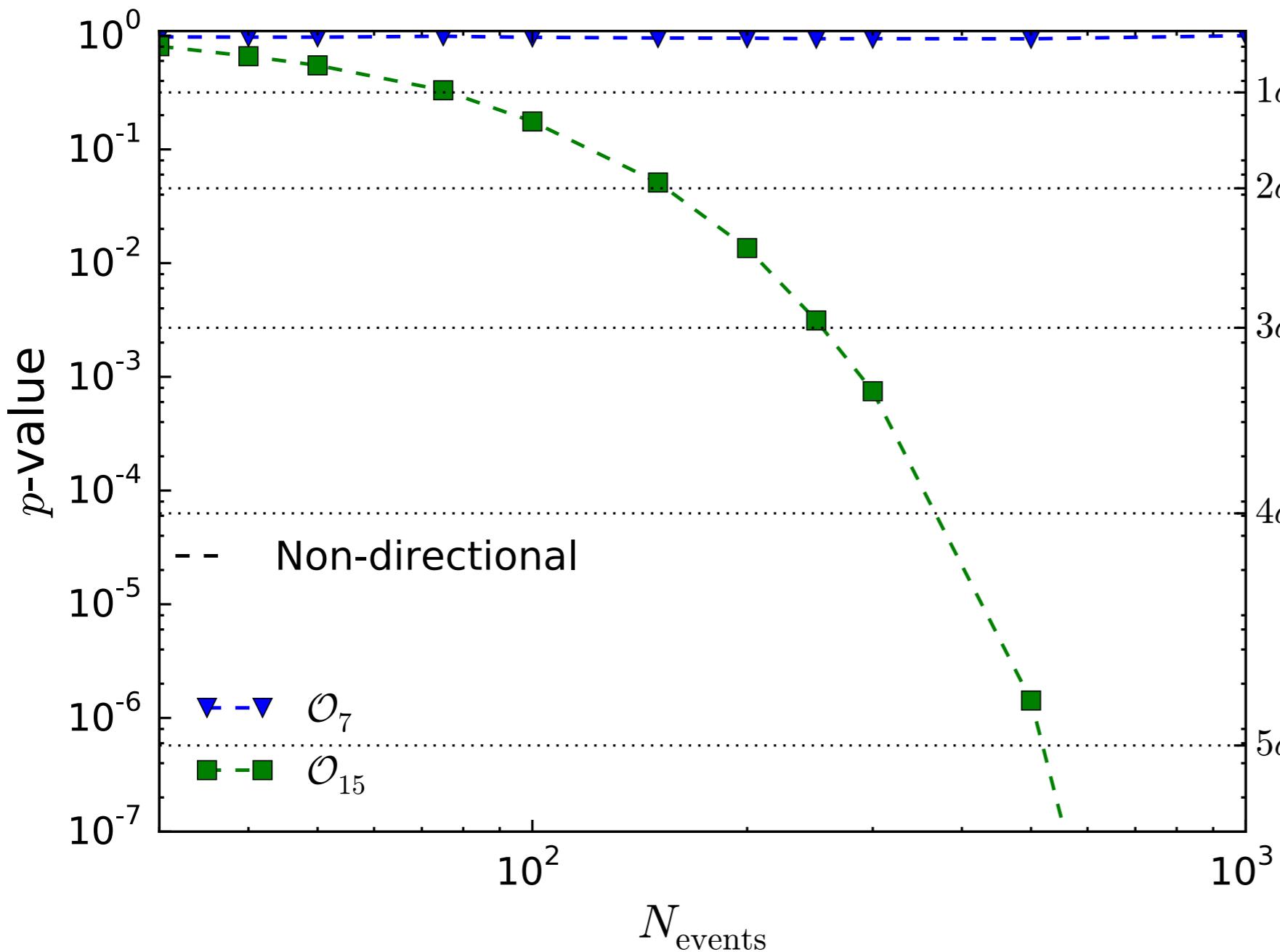
$$F_4 \sim 1$$

$$F_7 \sim v_\perp^2$$

$$F_{15} \sim q^2(q^2 + v_\perp^2)$$

# Distinguishing operators: Energy-only

*How many events are required to detect the effect of a ‘non-standard’ operator?*



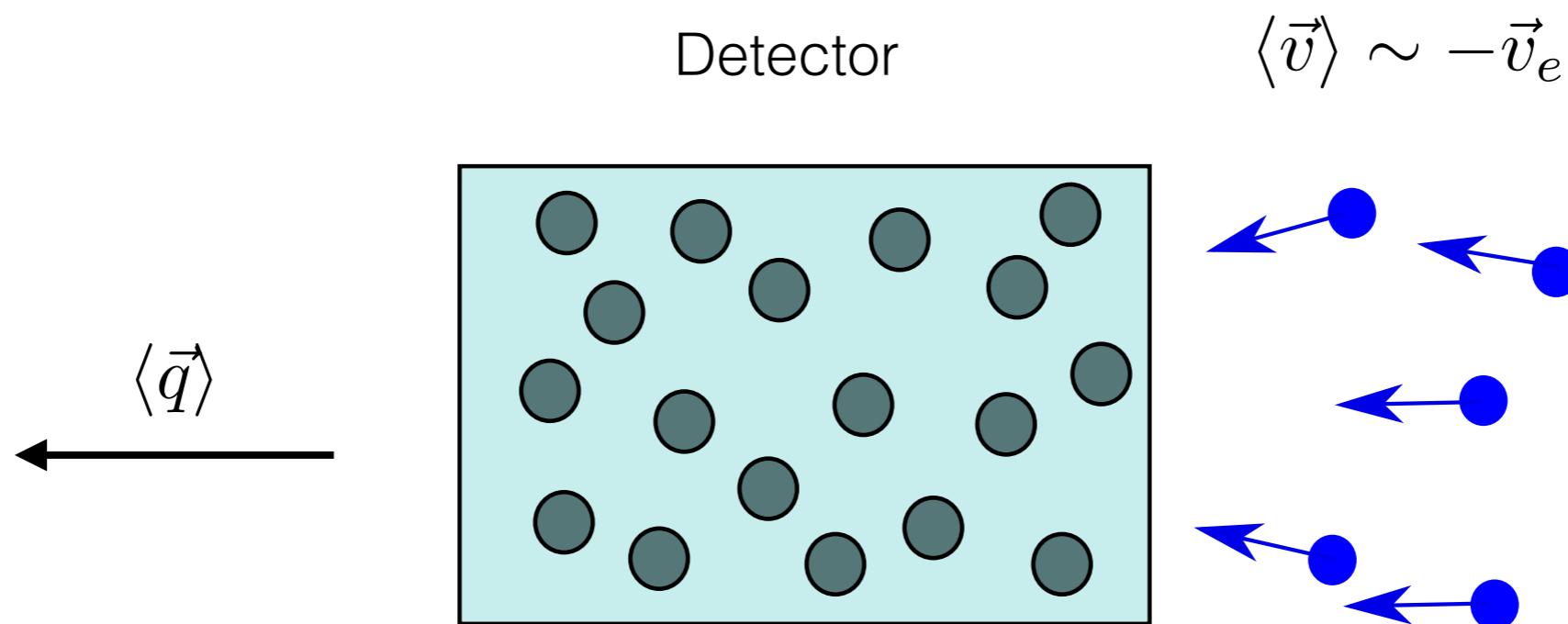
$F_4 \sim 1$   
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# Directional Detection

Different  $v$ -dependence could impact *directional* signal.

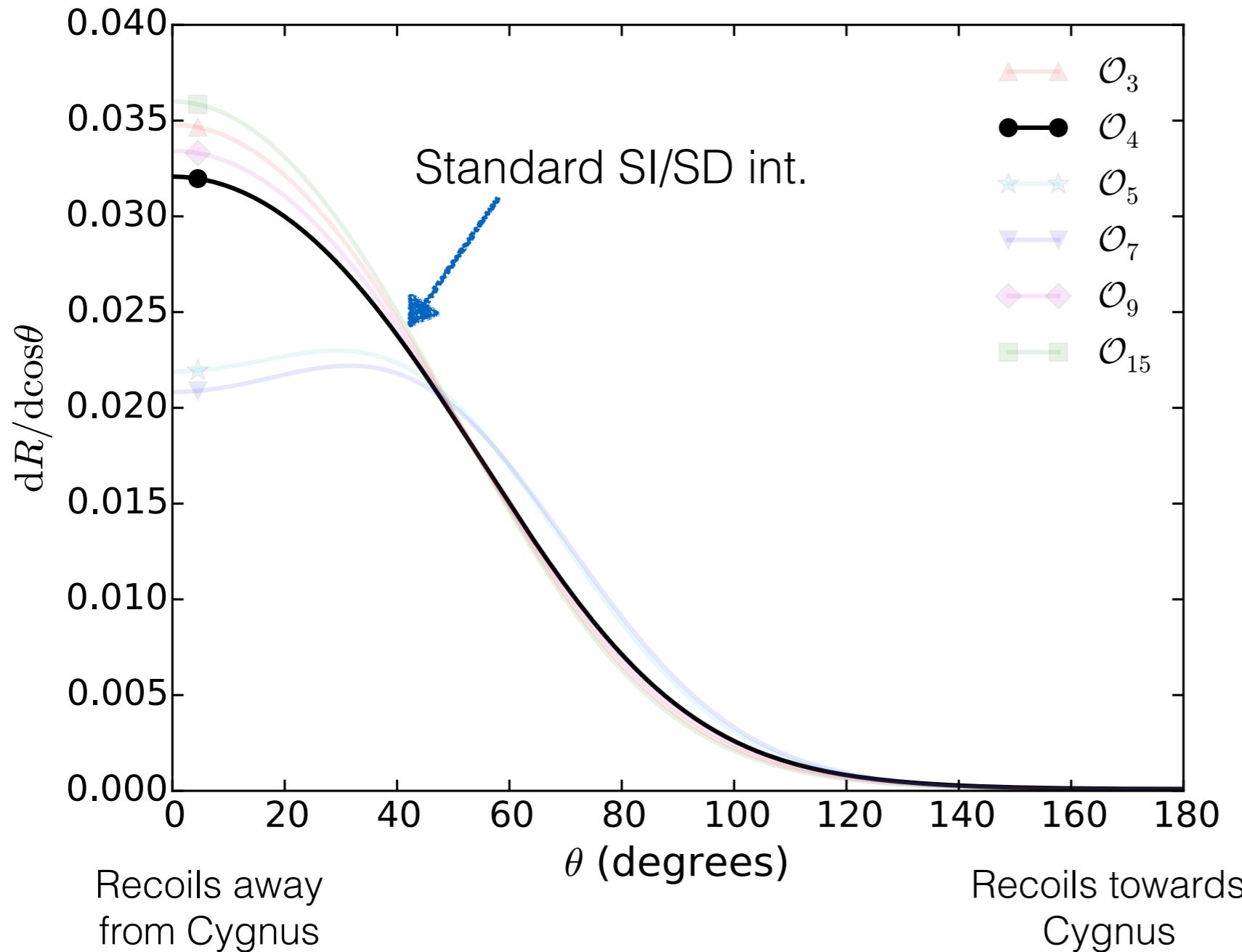
e.g. Drift-IIId [arXiv:1010.3027]

Mean recoil direction should point away from constellation Cygnus, due to Earth's motion.



Look at recoil rate, as a function of  $\theta$ , the angle between the recoil and the mean recoil direction.

# Standard directional spectrum



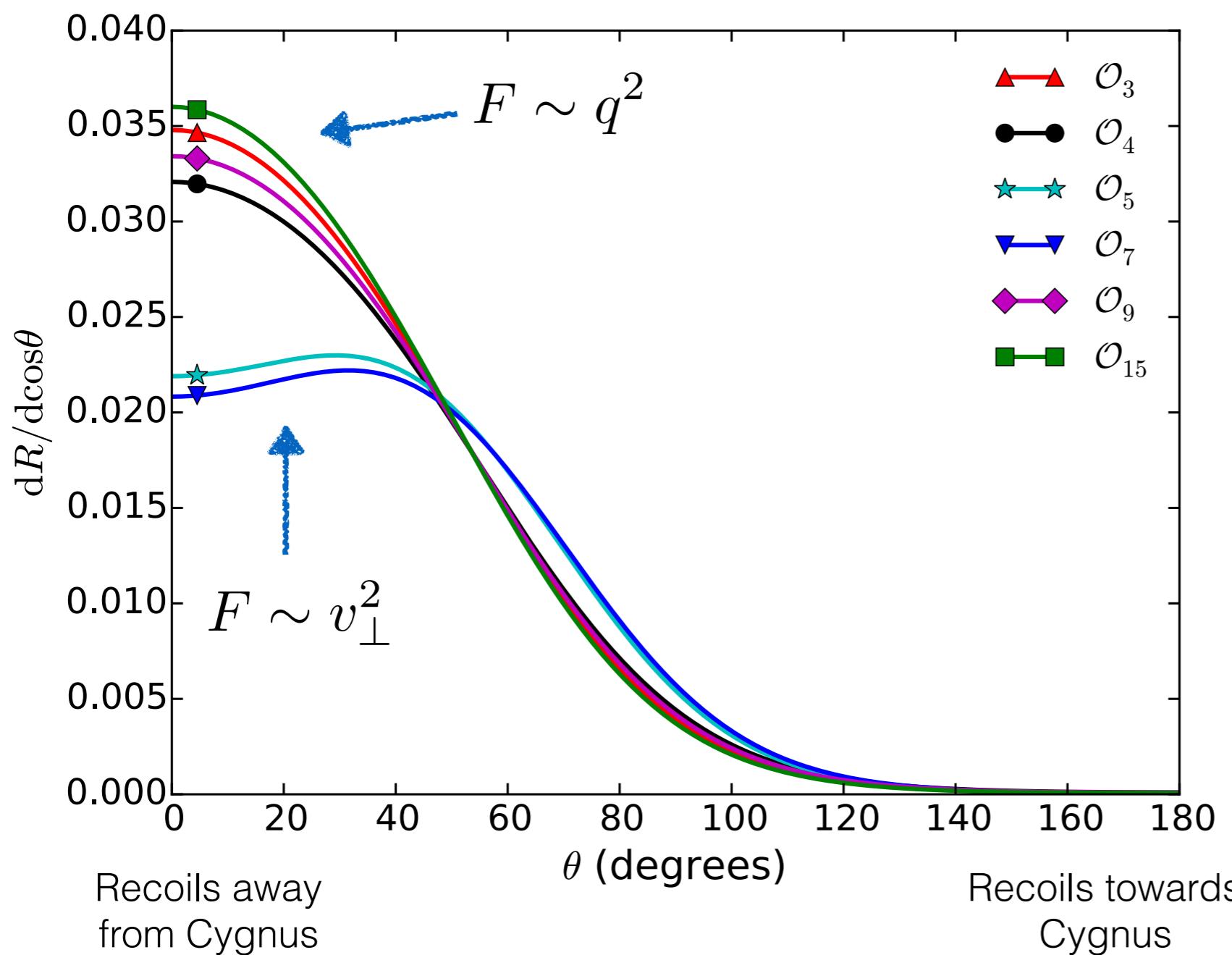
'Perfect' CF<sub>4</sub> detector

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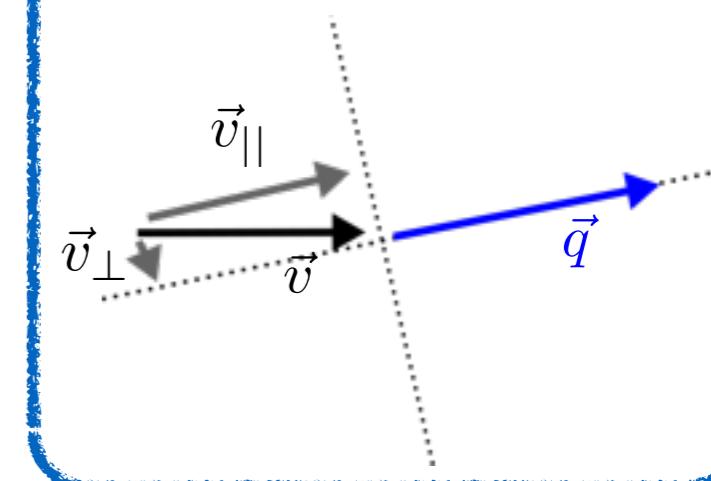
Input WIMP mass:  
 $m_\chi = 100$  GeV

SHM velocity distribution

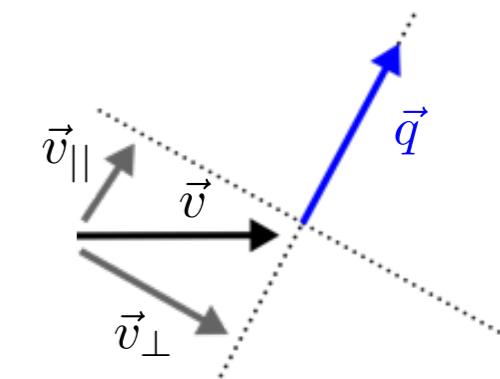
# NREFT directional spectrum



small  $\theta$ , small  $v_\perp$

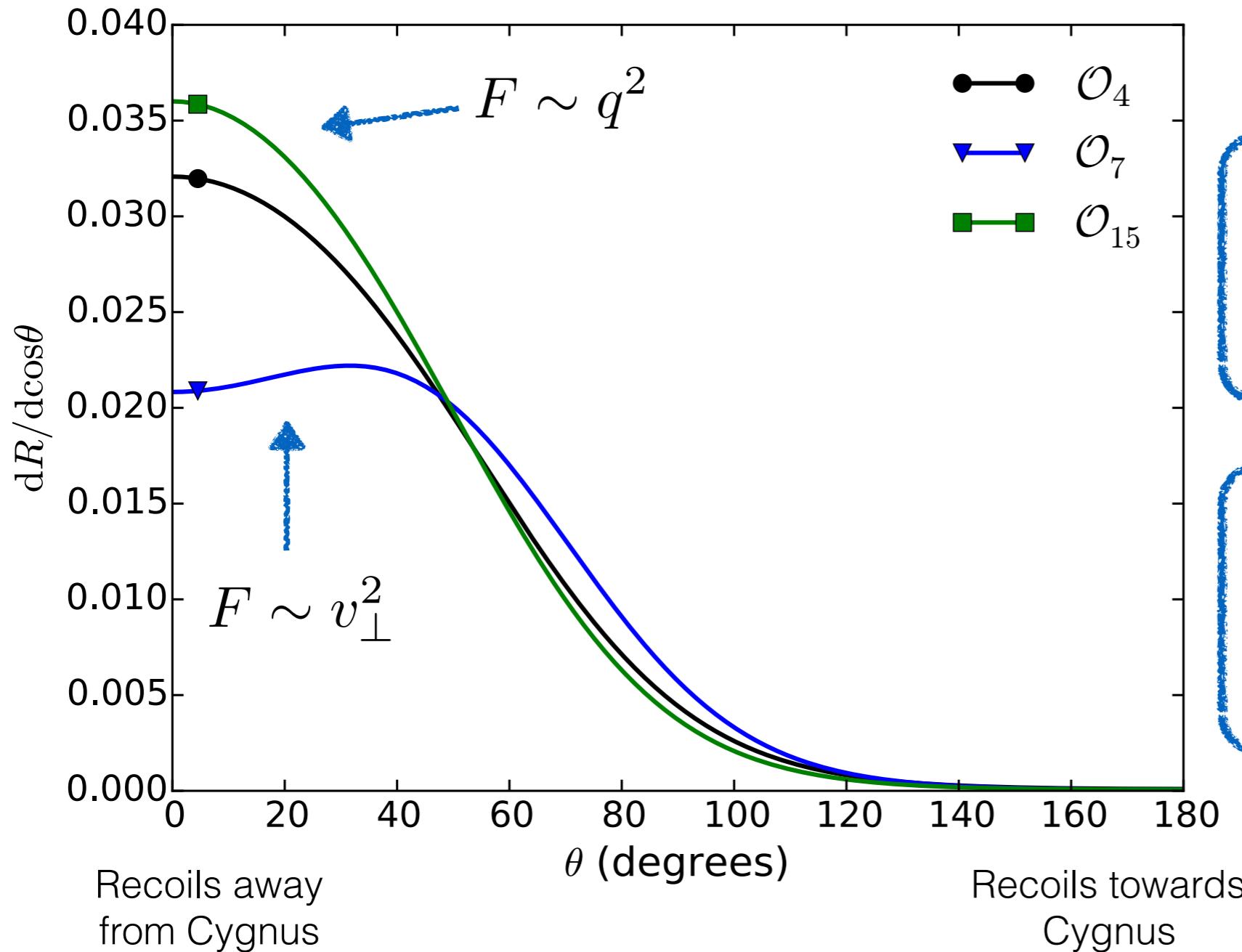


large  $\theta$ , large  $v_\perp$



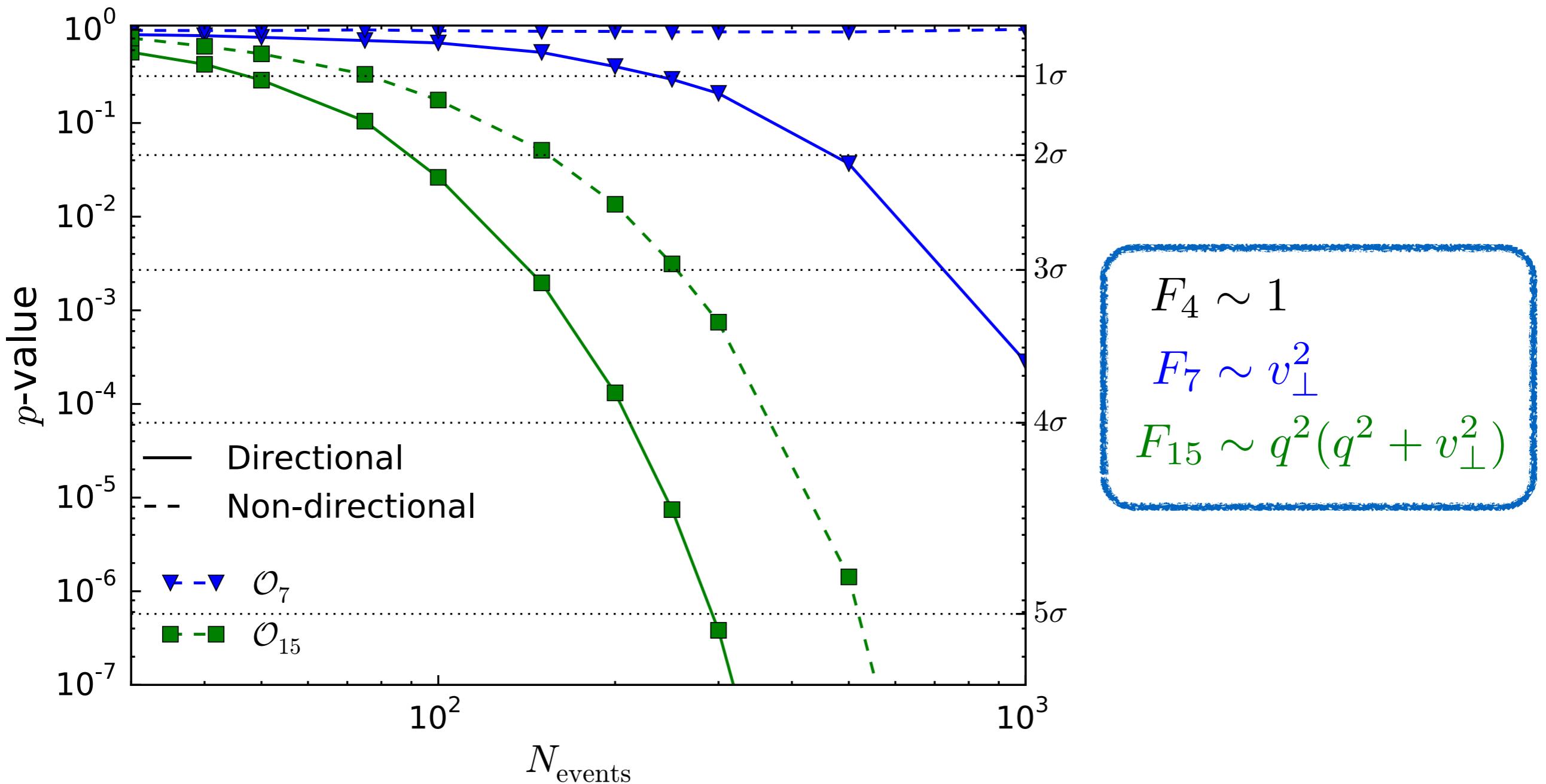
$$\begin{aligned} \text{Also note: } q &= 2\mu_{\chi N} \vec{v} \cdot \hat{q} \\ &= 2\mu_{\chi N} v \cos \theta \end{aligned}$$

# NREFT directional spectrum



# Distinguishing operators: Energy and direction

*How many events are required to detect the effect of a ‘non-standard’ operator?*



# Conclusions

Direct detection is a unique probe of the different possible interactions between DM and nucleons.

However, not all operators can be distinguished in an energy-only experiment. E.g.:

$$\mathcal{L}_1 = \bar{\chi}\chi\bar{N}N \longrightarrow F \sim v^0$$

$$\mathcal{L}_6 = \bar{\chi}\gamma^\mu\gamma^5\chi\bar{N}\gamma_\mu N \longrightarrow F \sim v_\perp^2$$

But, many operators have interesting *directional* signatures and directional sensitivity may allow us to detect the effects of ‘non-standard’ operators with only a few hundred events.

*Directional detection allows us to probe otherwise inaccessible particle physics of Dark Matter!*

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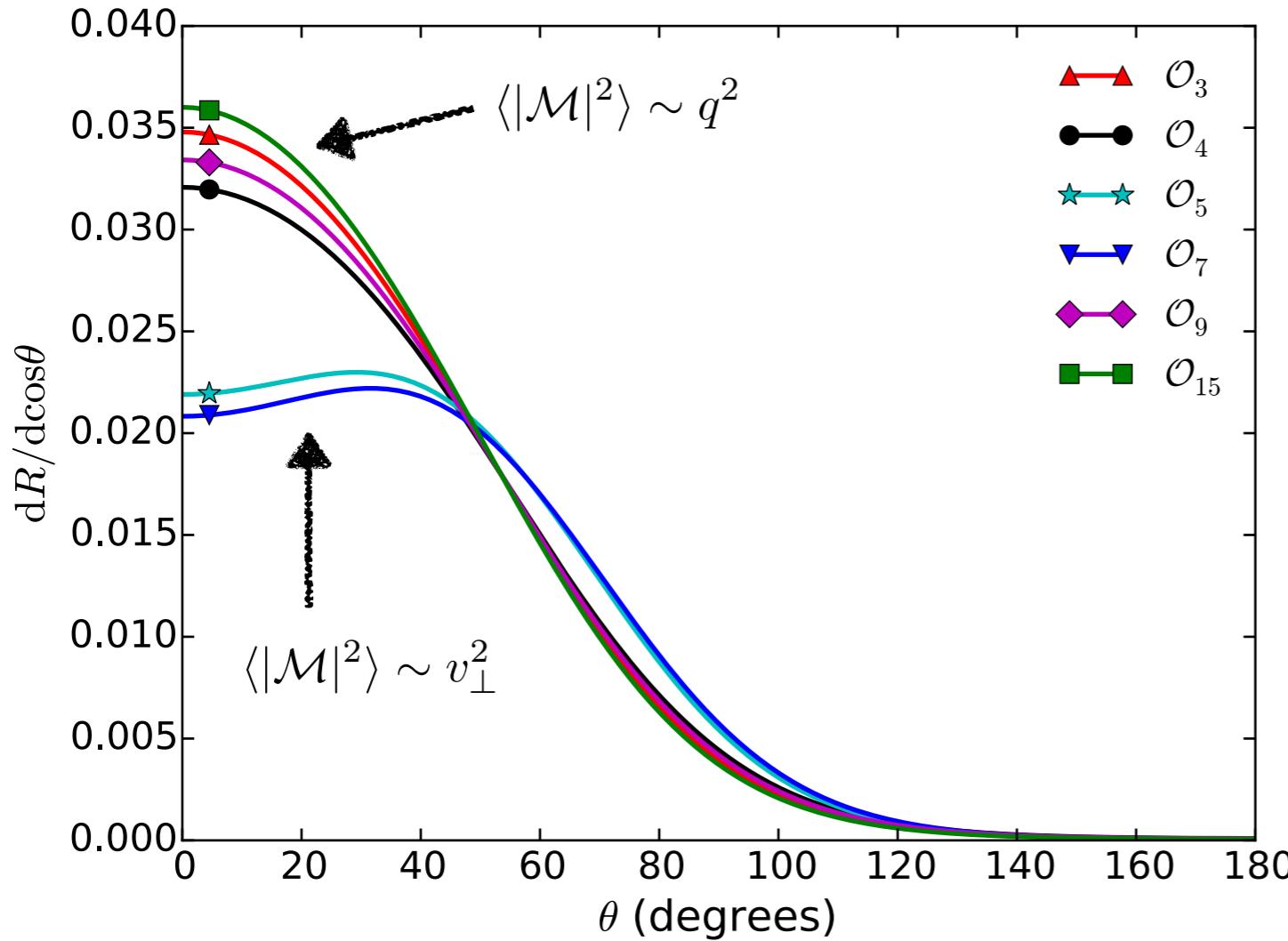
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*Directional detection allows us to probe otherwise inaccessible particle physics of Dark Matter!*

**Thank you**

# Backup Slides

# Directional Spectra



$$\langle |M|^2 \rangle \sim \begin{cases} 1 & : \mathcal{O}_1, \mathcal{O}_4, \\ v_\perp^2 & : \mathcal{O}_7, \mathcal{O}_8, \\ q^2 & : \mathcal{O}_9, \mathcal{O}_{10}, \mathcal{O}_{11}, \mathcal{O}_{12}, \\ v_\perp^2 q^2 & : \mathcal{O}_5, \mathcal{O}_{13}, \mathcal{O}_{14}, \\ q^4 & : \mathcal{O}_3, \mathcal{O}_6, \\ q^4(q^2 + v_\perp^2) & : \mathcal{O}_{15}. \end{cases}$$

Note:  $q = 2\mu_{\chi N} \vec{v} \cdot \hat{q}$   
 $= 2\mu_{\chi N} v \cos \theta$

Most isotropic:

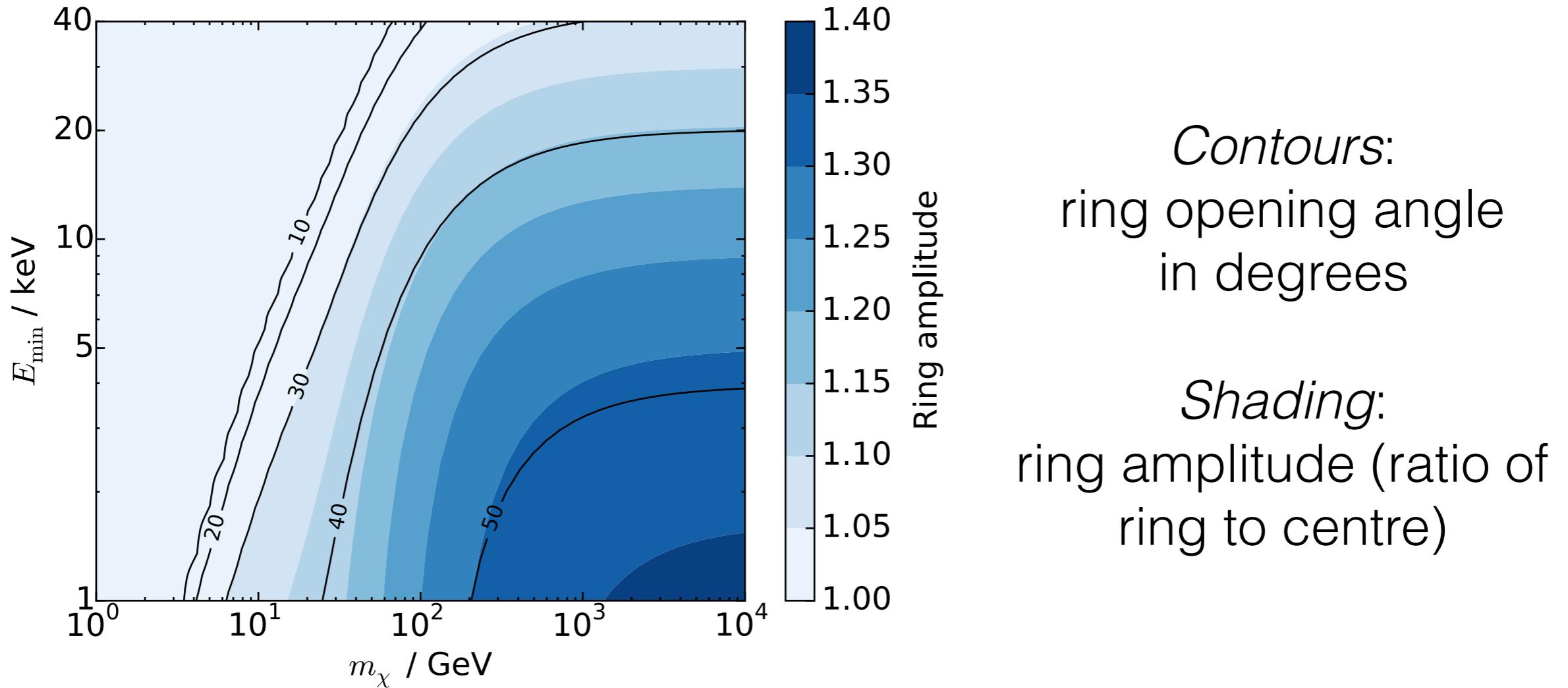
$$\mathcal{O}_7 = \vec{S}_n \cdot \vec{v}_\perp$$

Least isotropic:

$$\mathcal{O}_{15} = (\vec{S}_\chi \cdot \frac{\vec{q}}{m_n}) ((\vec{S}_n \times \vec{v}_\perp) \cdot \frac{\vec{q}}{m_n})$$

# A (new) ring-like feature

Operators with  $\langle |\mathcal{M}|^2 \rangle \sim (v_\perp)^2$  lead to a ‘ring’ in the directional rate.



A ring in the standard rate has been previously studied [Bozorgnia et al. - 1111.6361], but *this* ring occurs for lower WIMP masses and higher threshold energies.

# Statistical tests

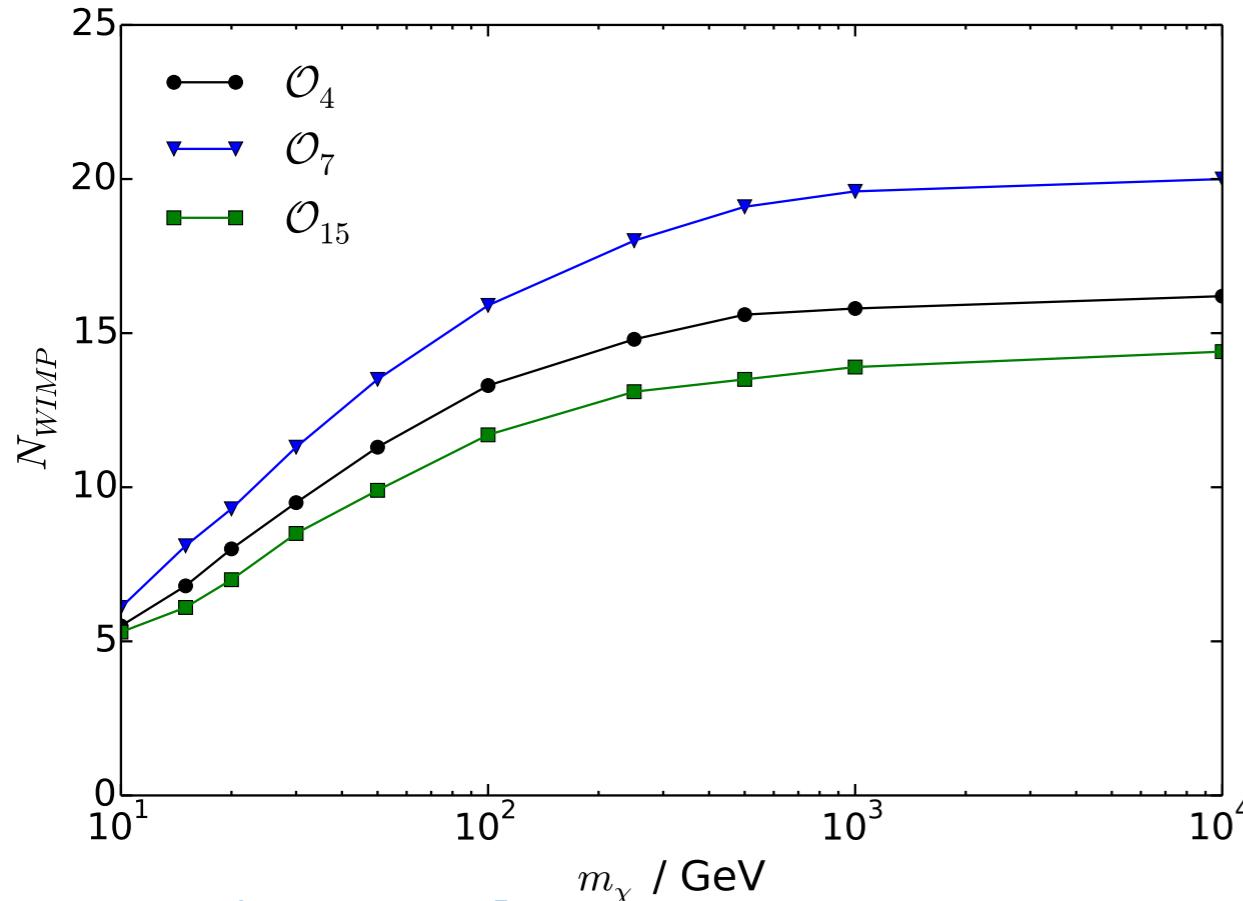
$$F_{4,4} \sim 1$$

$$F_{7,7} \sim v_{\perp}^2$$

$$F_{15,15} \sim q^4(q^2 + v_{\perp}^2)$$

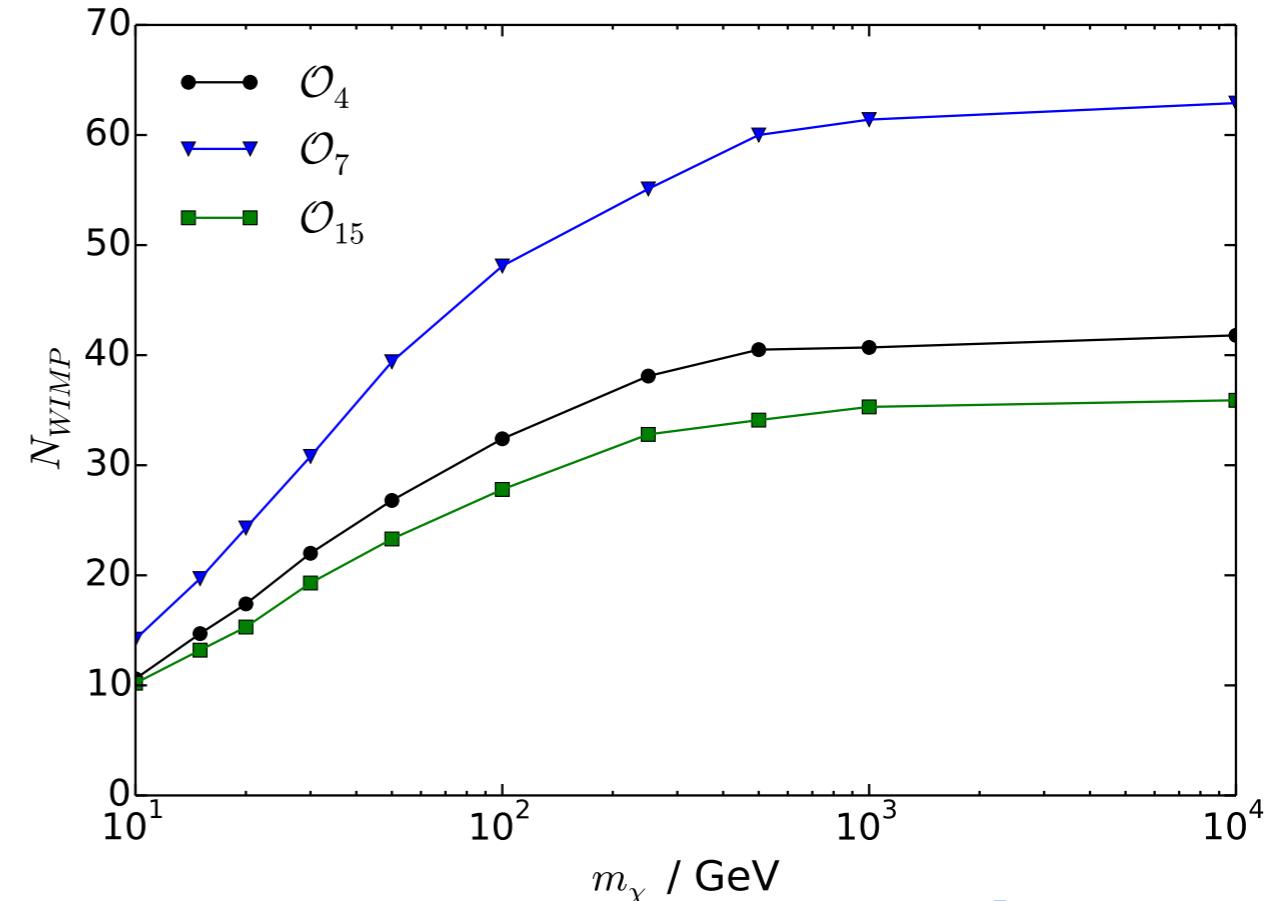
Calculate the number of signal events required to...

...reject isotropy...



[astro-ph/0408047]

...confirm the median recoil dir...



[1002.2717]

...at the  $2\sigma$  level in 95% of experiment.

# Distinguishing operators

Generate data assuming an NREFT operator (  $\mathcal{O}_7$  or  $\mathcal{O}_{15}$  ).

Assume data is a combination of standard SI/SD interaction and non-standard NREFT operator. Fit to data with two free parameters  $m_\chi$  and  $A$ .

$A$  : fraction of events which are due to non-standard NREFT interaction.

Perform likelihood ratio test (in 10000 pseudo-experiments) to determine the significance with which we can reject SD-only interactions:

Null hypothesis,  $\mathbf{H}_0$ : all events are due to SD interactions,  $A = 0$

Alt. hypothesis,  $\mathbf{H}_1$ : there is some contribution from NREFT ops,  $A \neq 0$

# Consequences for relativistic theories

Many ‘dictionaries’ are available which allow us to translate from relativistic interactions to NREFT interactions

[e.g. 1211.2818, 1307.5955, 1505.03117]

$$\begin{aligned}\mathcal{L}_1 = \bar{\chi}\chi\bar{n}n &\quad \Rightarrow \quad \langle \mathcal{L}_1 \rangle = 4m_\chi m_n \mathcal{O}_1 \\ &\quad \rightarrow \quad \langle |\mathcal{M}|^2 \rangle \sim F_M(q^2)\end{aligned}$$

$$\begin{aligned}\mathcal{L}_6 = \bar{\chi}\gamma^\mu\gamma^5\chi\bar{n}\gamma_\mu n &\quad \Rightarrow \quad \langle \mathcal{L}_6 \rangle = 8m_\chi(m_n \mathcal{O}_8 + \mathcal{O}_9) \\ &\quad \rightarrow \quad \langle |\mathcal{M}|^2 \rangle \sim v_\perp^2 F_M(q^2)\end{aligned}$$

*These two relativistic operators cannot be distinguished without directional detection.*

# Open issues

- We have assumed an *ideal* detector - lower limits on the event numbers (need to be convolved with detector effects...)
- Different signatures possible for different target materials - see (very) recent paper by Catena [\[1505.06441\]](#)
- Astrophysical uncertainties are expected to be comparable with particle physics uncertainties
  - inability to distinguish different operators depends on SHM-type distribution (may be different for sharp stream-like distributions)
- May be possible to distinguish operators using other methods
  - measuring annual modulation [\[1504.06772\]](#)

In the future, it would be interesting to examine astrophysical uncertainties in detail, and to compare different approaches to distinguishing NREFT operators.